# Nonlinear Self-organization Dynamics of a Metabolic Process of the Krebs Cycle

Valeriy Grytsay<sup>1</sup>, Iryna Musatenko<sup>2</sup>

**Abstract.** The present work continues studies of the mathematical model of a metabolic process of the Krebs cycle. We study the dependence of its cyclicity on the cell respiration intensity determined by the formation level of carbon dioxide. We constructed the phase-parametric characteristic of the consumption of a substrate by a cell depending on the intensity of the metabolic process of formation of the final product of the oxidation. The scenarios of all possible oscillatory modes of the system are constructed and studied. The bifurcations with period doubling and with formation of chaotic modes are found. Their attractors are constructed. The full spectra of indices and divergencies for the obtained modes, the values of KS-entropies, horizons of predictability, and Lyapunov dimensions of strange attractors are calculated. Some conclusions about the structural-functional connections of the cycle of tricarboxylic acids and their influence on the stability of the metabolic process in a cell are presented.

**Keywords:** Krebs cycle, metabolic process, self-organization, strange attractor, bifurcation, Feigenbaum scenario.

# 1 Introduction

One of the possible problems of synergetics is the study of the internal dynamics of metabolic processes in cells. Its solution allows one to find the structural-functional connections defining the self-organization of these processes and to answer the question how the catalyzed enzymatic reactions create the internal space-time ordering of the cell life.

The most general metabolic process in cells is the cycle of tricarboxylic acids [1]. This is the key stage of the respiration of all cells. In its course, the diand tricarbon compounds, which are formed as intermediate products in the transformation of carbohydrates, fats, and proteins, are transformed up to CO<sub>2</sub>. In this case, the released hydrogen is oxidized further up to water, by taking the direct participation in the synthesis of ATP, being the universal energy source.

Studies of the functioning of the Krebs cycle were carried out both experimentally and theretically in [2-10].

In the study of the given process, we use the mathematical model of the growh of cells *Candida utilis* on ethanol, which was developed by Professor

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Bogolyubov Institute for Theoretical Physics, 14b, Metrolohichna Str., Kyiv 03680 (E-mail: vgrytsay@bitp.kiev.ua)

<sup>&</sup>lt;sup>2</sup> Taras Shevchenko National University of Kyiv, Faculty of Cybernetics, Department of Computational Mathematics, 64, Volodymyrska Str., Kyiv, Ukraine (E-mail: ivmusatenko@gmail.com)

V.P. Gachok [11, 12]. With the help of this model, the unstable modes in the cultivation of cells observed in experiments were considered. The kinetic curves of the chaotic dynamics obtained with the help of computational experiments were in agreement with experimental data [13].

Then the given model was modified and refined in [14] due to the account for the influence of the CO<sub>2</sub> level on the respiration intensity. With the help of the model, the structural-functional connections of the metabolic process in a cell, which cause the appearance of complicated oscillations in the metabolic process, were investigated. It was concluded that the given oscillations arise on the level of redox reactions of the Krebs cycle, reflect the cyclicity of the process, and characterize the self-organization in a cell. The fractality of the dynamics of oscillations of the Krebs cycle was studied as well.

The analogous oscillatory modes were observed in the processes of photosynthesis and glycolysis, variations of the calcium concentration in a cell, oscillations in heart muscle, and other biochemical processes [15-19].

#### 2 Mathematical Model

The general scheme of the process is presented in Fig. 1. According to it with regard for the mass balance, we have constructed the mathematical model given by Eqs. (1) - (19).

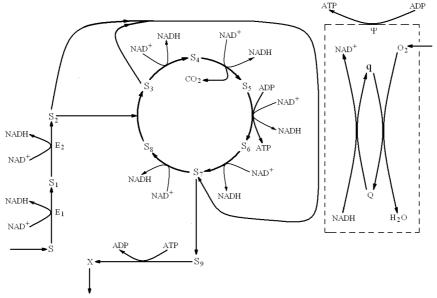


Fig. 1. General scheme of the metabolic process of growth of cells *Candida utilis* on ethanol.

$$\frac{dS}{dt} = S_0 \frac{K}{K + S + \gamma \psi} - k_1 V(E_1) \frac{N}{K_1 + N} V(S) - \alpha_1 S,$$
(1)

$$\frac{dS_1}{dt} = k_1 V(E_1) \frac{N}{K_1 + N} V(S) - k_2 V(E_2) \frac{N}{K_1 + N} V(S_1), \tag{2}$$

$$\frac{dS_2}{dt} = k_2 V(E_2) \frac{N}{K_1 + N} V(S_1) - k_3 V(S_2^2) V(S_3) - k_4 V(S_2) V(S_8), \tag{3}$$

$$\frac{dS_3}{dt} = k_4 V(S_2) V(S_8) - k_5 V(N^2) V(S_3^2) - k_3 V(S_2^2) V(S_3), \tag{4}$$

$$\frac{dS_4}{dt} = k_5 V(N^2) V(S_3^2) - k_7 V(N) V(S_4) - k_8 V(N) V(S_4), \tag{5}$$

$$\frac{dS_5}{dt} = k_7 V(N) V(S_4) - 2k_9 V(L_1 - T) V(S_5), \tag{6}$$

$$\frac{dS_6}{dt} = 2k_9V(L_1 - T)V(S_5) - k_{10}V(N)\frac{S_6^2}{S_6^2 + 1 + M_1S_8},\tag{7}$$

$$\frac{dS_7}{dt} = k_{10}V(N)\frac{S_6^2}{S_6^2 + 1 + M_1S_8} - k_{11}V(N)V(S_7) - K_{12}\frac{S_7^2}{S_7^2 + 1 + M_2S_9}V(\psi^2) + k_3V(S_2^2)V(S_3),$$
(8)

$$\frac{dS_8}{dt} = k_{11}V(N)V(S_7) - k_4V(S_2)V(S_8) + k_6V(T^2)\frac{S^2}{S^2 + \beta_1} \cdot \frac{N_1}{N_1 + (S_5 + S_7)^2},$$
 (9)

$$\frac{dS_9}{dt} = k_{12} \frac{S_7^2}{S_7^2 + 1 + M_2 S_9} V(\psi^2) - k_{14} \frac{XTS_9}{(\mu_1 + T)[(\mu_2 + S_9 + X + M_3(1 + \mu_3 \psi)]S}, (10)$$

$$\frac{dX}{dt} = k_{14} \frac{XTS_9}{(\mu_1 + T)[(\mu_2 + S_9 + X + M_3(1 + \mu_3 \psi)]S} - \alpha_2 X,\tag{11}$$

$$\frac{dQ}{dt} = -k_{15}V(Q)V(L_2 - N) + 4k_{16}V(L_3 - Q)V(O_2)\frac{1}{1 + \gamma_1 \psi^2},$$
(12)

$$\frac{dO_2}{dt} = O_{2_0} \frac{K_2}{K_2 + O_2} - k_{16}(L_3 - Q)V(O_2) \frac{1}{1 + \gamma_1 \psi} - k_8 V(N)V(S_4) - \alpha_3 O_2, \tag{13}$$

$$\frac{dN}{dt} = -k_7 V(N) V(S_4) - k_{10} V(N) \frac{S_6^2}{S_6^2 + 1 + M_1 S_8} - k_{11} V(N) V(S_7) - k_5 V(N^2) V(S_3^2) + k_{15} V(Q) V(L_2 - N) - k_2 V(E_2) \frac{N}{K_1 + N} V(S_1) - k_1 V(E_1) \frac{N}{K_1 + N} V(S),$$
(14)

$$\frac{dT}{dt} = k_{17}V(L_1 - T)V(\psi^2) + k_9V(L - T)V(S_3) - \alpha_4 T - \\
- k_{18}k_6V(T^2) \frac{S^2}{S^2 + \beta_1} \cdot \frac{N_1}{N_1 + (S_5 + S_7)^2} - \\
- k_{19}k_{14} \frac{XTS_9}{(\mu_1 + T)[\mu_2 + S_9 + X + M_3(1 + \mu_3 \psi)S]},$$
(15)

$$\frac{d\psi}{dt} = 4k_{15}V(Q)V(L_2 - N) + 4k_{17}V((L_1 - T)V(\psi^2) - \frac{S_7^2}{S_7^2 + 1 + M_2S_9}V(\psi^2) - \alpha\psi,$$
(16)

$$\frac{dE_1}{dt} = E_{1_0} \frac{S^2}{\beta_2 + S^2} \frac{N_2}{N_2 + S_1} - n_1 V(E_1) \frac{N}{K_1 + N} V(S) - \alpha_5 E_1, \tag{17}$$

$$\frac{dE_2}{dt} = E_{2_0} \frac{S_1^2}{\beta_3 + S_1^2} \frac{N_3}{N_3 + S_2} - n_2 V(E_2) \frac{N}{K_1 + N} V(S_1) - \alpha_6 E_2, \tag{18}$$

$$\frac{dC}{dt} = k_8 V(N)V(S_4) - \alpha_7 C. \tag{19}$$

where V(X) = X/(1+X) is the function that describes the adsorption of the enzyme in the region of a local coupling. The variables of the system are dimensionless [11, 12].

The internal parameters of the system are as follows:

The external parameters determining the flow-type conditions are chosen as  $S_0 = 0.05055$ ;  $O_{2_0} = 0.06$ ;  $\alpha = 0.002$ ;  $\alpha_1 = 0.02$ ;  $\alpha_2 = 0.004$ ;  $\alpha_3 = 0.01$ ;  $\alpha_4 = 0.01$ ;  $\alpha_5 = 0.01$ ;  $\alpha_6 = 0.01$ ;  $\alpha_7 = 0.0001$ .

The model covers the processes of substrate-enzymatic oxidation of ethanol to acetate, cycle involving tri- and dicarboxylic acids, glyoxylate cycle, and respiratory chain.

The incoming ethanol S is oxidized by the alcohol dehydrogenase enzyme  $E_1$  to acetaldehyde  $S_1$  (1) and then by the acetal dehydrogenase enzyme  $E_2$  to acetate  $S_2$  (2), (3). The formed acetate can participate in the cell metabolism and can be exchanged with the environment. The model accounts for this situation by the change of acetate by acetyl-CoA. On the first stage of the Krebs

cycle due to the citrate synthase reaction, acetyl-CoA jointly with oxalacetate  $S_8$  formed in the Krebs cycle create citrate  $S_3$  (4). Then substances  $S_4$  -  $S_8$ are created successively on stages (5)-(9). In the model, the Krebs cycle is represented by only those substrates that participate in the reduction of NADH and the phosphorylation  $ADT \rightarrow ATP$ . Acetyl- CoA passes along the chain to malate represented in the model as intramitochondrial  $S_7$  (8) and cytosolic  $S_9$ (10) ones. Malate can be also synthesized in another way related to the activity of two enzymes: isocitrate lyase and malate synthetase. The former catalyzes the splitting of isocitrate to succinate, and the latter catalyzes the condensation of acetyl- CoA with glyoxylate and the formation of malate. This glyoxylate-linked way is shown in Fig. 1 as an enzymatic reaction with the consumption of  $S_2$ and  $S_3$  and the formation of  $S_7$ . The parameter  $k_3$  controls the activity of the активность glyoxylate-linked way (3), (4), (8). The yield of  $S_7$  into cytosol is controlled by its concentration, which can increase due to  $S_9$ , by causing the inhibition of its transport with the participation of protons of mitochondrial membrane.

The formed malate  $S_9$  is used by a cell for its growth, namely for the biosynthesis of protein X (11). The energy consumption of the given process is supported by the process  $ATP \rightarrow ADP$ . The presence of ethanol in the external solution causes the "ageing" of external membranes of cells, which leads to the inhibition of this process. The inhibition of the process also happens due to the enhanced level of the kinetic membrane potential  $\psi$ . The parameter  $\mu_0$  is related to the lysis and the washout of cells.

In the model, the respiratory chain of a cell is represented in two forms: oxidized, Q, (12) and reduced, q, ones. They obey the integral of motion  $Q(t)+q(t)=L_3$ .

A change of the concentration of oxygen in the respiratory chain is determined by Eq. (13).

The activity of the respiratory chain is affected by the level of NADH (14). Its high concentration leads to the enhanced endogenic respiration in the reducing process in the respiratory chain (parameter  $k_{15}$ ). The accumulation of

NADH occurs as a result of the reduction of  $NAD^+$  at the transformation of ethanol and in the Krebs cycle. These variables obey the integral of motion  $NAD^+(t) + NADH(t) = L_2$ .

In the respiratory chain and the Krebs cycle, the substrate-linked phosphorylation of ADP with the formation of ATP (15) is also realized. The energy consumption due to the process  $ATP \rightarrow ADP$  induces the biosynthesis of components of the Krebs cycle (parameter  $k_{18}$ ) and the growth of cells on the substrate (parameter  $k_{19}$ ). For these variables, the integral of motion  $ATP(t) + ADP(t) = L_1$  holds. Thus, the level of ATP produced in the redox

processes in the respiratory chain  $ADP \rightarrow ATP$  determines the intensity of the Krebs cycle and the biosynthesis of protein.

In the respiratory chain, the kinetic membrane potential  $\psi$  (16) is created under the running of reducing processes  $Q \to q$ . It is consumed at the substrate-linked phosphorylation  $ADP \to ATP$  in the respiratory chain and the Krebs cycle. Its enhanced level inhibits the biosynthesis of protein and process of reduction of the respiratory chain.

Equations (17) and (18) describe the activity of enzymes  $E_1$  and  $E_2$ , respectively. We consider their biosynthesis ( $E_{1_0}$  and  $E_{2_0}$ ), the inactivation in the course of the enzymatic reaction ( $n_1$  and  $n_2$ ), and all possible irreversible inactivations ( $\alpha_5$  and  $\alpha_6$ ).

Equation (19) is related to the formation of carbon dioxide. Its removal from the solution into the environment ( $\alpha_7$ ) is taken into account. Carbon dioxide is produced in the Krebs cycle (5). In addition, it squeezes out oxygen from the solution (13), by decreasing the activity of the respiratory chain.

The study of solutions of the given mathematical model (1)-(19) was performed with the help of the theory of nonlinear differential equations [20, 21] and the methods of mathematical modeling of biochemical systems applied and developed by the authors. in [22-38].

## 3 The results of Studies

For one cycle, there occurs the full oxidation of a molecule of acetyl- *CoA* up to malate and the formation of a new molecule of acetyl- *CoA* at the input. In such a way, the continuous process of functioning of the Krebs cycle is running. This process has the autooscillatory character.

The studies of the model with the help of computational experiments showed that if system's parameters vary, the appearance of autooscillations with various frequencies, as well as chaotic oscillations, becomes possible. Oscillations with the same frequency will occur in all components of the given metabolic process. In the present work, we will study the dependence of autooscillations of the system on the parameter  $k_8$ , which determines the level of formation of  $\mathrm{CO}_2$  in the cycle of tricarboxylic acid.

The different types of obtained autooscillatory modes are studied with the help of the construction of phase-parametric diagrams. The abscissa axis shows the values of parameter  $k_8$ , and the axis of ordinates gives the values of chosen variable  $E_1(t)$ , for example. Moreover, we used the method of cutting. In the phase space of trajectories of the system, we place the cutting plane  $S_2=0.8$ . Such a choice is explained by the symmetry of oscillations of acetate relative to this plane in a lot of earlier calculated modes. For every given value of  $k_8$ , we observe the intersection of this plane by the trajectory in a single direction, when it approaches the attractor. The value of  $E_1(t)$  is put onto the phase-

parametric diagram. In the case where a multiple periodic limiting cycle arises, a number of points can be observed on the plane, and they will be the same in the period. If the deterministic chaos arises, the points of the intersection of the plane by the oscillating trajectory will be positioned chaotically.

In Fig. 2,a-d, we show the phase-parametric diagrams for the variable  $E_1(t)$  versus the parameter  $k_8$  changing in the appropriate intervals.

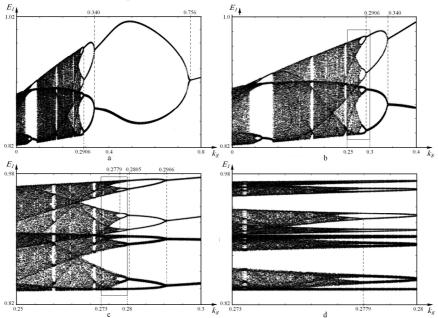


Fig. 2. Phase-parametric diagram for the variable  $E_1(t)$ : a -  $k_8 \in (0.,0.8)$ ; b -  $k_8 \in (0.,0.4)$ ; c -  $k_8 \in (0.25,0.3)$ ; d -  $k_8 \in (0.273,0.28)$ .

As the parameter  $k_8$  decreases, there occurs the subsequent doubling of the multilicity of the autoperiodic process. Such a sequence of the appearance of bifurcations creates a cascade of bifurcations, namely the Feigenbaum sequence [39]. After the multiple doubling of a period, the modes of aperiodic oscillations are eventually observed in the system. In other words, a chaos arises. As the parameter  $k_8$  decreases further, we see the appearance of the windows of periodicity on the phase-parametric diagrams. The deterministic chaos is destroyed, and the periodic and quasiperiodic modes are established. The trajectory of a strange attractor in the chaotic mode is tightened to a regular attractor of the autoperiodic mode. We observe the self-organization in the system. Then the windows of periodicity are destroyed, and the chaotic modes arise again. Moreover, the transitions "order—chaos" and "chaos—order" happen. There occurs the adaptation of the metabolic process to varying conditions.

It is seen from the presented figures that, as the scale decreases, every subsequent phase-parametric diagram with doubling of a cycle and its windows

of periodicity are identicat to those of the previous diagram, as the scale decreases. The given sequence of bifurcations has a self-similar fractal structure.

In Figs. 3,e-f and 4, we present the examples of the projections of phase portraits for some values of parameter  $k_8$ , according to the phase-parametric diagram in Fig. 2.

In Fig. 5, we show the constructed kinetic curves for a strange attractor formed at  $k_8 = 0.12$ .

These figures indicate a variation of the dynamics of a metabolic process of the Krebs cycle, which depends on the intensity of formation of the final oxidation product, CO<sub>2</sub>.

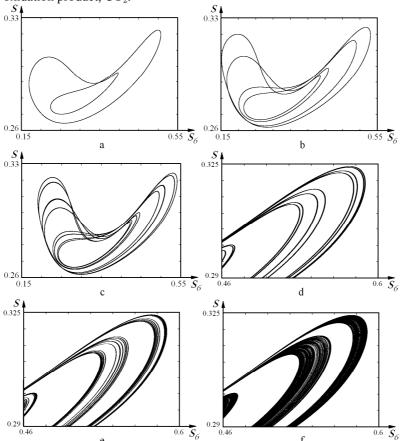
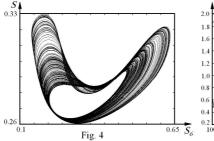


Fig. 3. Projections of system's phase portraits: a – regular attractor  $2 \cdot 2^1$ ,  $k_8 = 0.5$ ; b – regular attractor  $2 \cdot 2^2$ ,  $k_8 = 0.3$ ; c - regular attractor  $2 \cdot 2^4$ ,  $k_8 = 0.28$ ; d - regular attractor  $2 \cdot 2^8$ ,  $k_8 = 0.278$ ; e - regular attractor  $2 \cdot 2^{16}$ ,  $k_8 = 0.277$ ; f - strange attractor  $2 \cdot 2^x$ ,  $k_8 = 0.275$ .



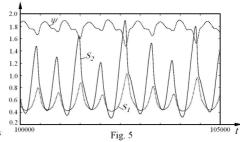


Fig. 4. Projection of the phase portrait of the strange attractor  $2 \cdot 2^x$  for  $k_8 = 0.12$ .

Fig. 5. Kinetic curve for the components  $S_1$ ,  $S_2$ , and  $\psi$  of the Krebs cycle in the mode of the strange attractor  $2 \cdot 2^x$  for  $k_8 = 0.12$ .

In order to uniquely identify the type of obrained attractors and to determine their stability, we calculated the full spectra of Lyapunov indices and

their sum  $\Lambda = \sum_{j=1}^{19} \lambda_j$  for the chosen points. The calculation was carried out by

Benettin's algorithm with the orthogonalization of the vectors of perturbations by the Gram--Schmidt method [21].

The calculation of Lyapunov indices from this multidimensional system on a personal computer meets certain difficulties. The mathematical model of the given biochemical system contains many variables and parameters. The limitations in the solution of such problems arise due to the insufficient random-access memory of a computer in the processing of the  $n \times n$  matrix of small perturbations. In addition, any inaccuracy on the stage of programming will essentially affect the redefinition of the vectors of perturbations, their orthogonalization, and, as a consequence, the result of calculations. Nevertheless, we solved the problem and obtained certain results. Below for the sake of comparison, we present the spectra of Lyapunov indices for some modes of the system. For brevity without any loss of information, we give the values of indices up to the fourth decimal point.

The ratios of the values of Lyapunov indices  $\lambda_1 > \lambda_2 > \lambda_3 > ... > \lambda_{19}$  serve as the criterion of the validity of calculations. For a regular attractor, we have obligatorily  $\lambda_1 \approx 0$ . The remaining indices can be also  $\approx 0$  in some cases. In some other cases, they are negative. The zero value of the first Lyapunov index testifies to the presence of a stable limiting cycle.

For a strange attractor, at least one Lyapunov index must be positive. After it, the zero index follows. The next indices are negative. The presence of negative indices means the contraction of system's phase space in the corresponding directions, whereas the positive indices indicate the dispersion of trajectories. Therefore, there occurs the mixing of trajectories in narrow places

For  $k_8=0.01$ , the strange attractor  $2\cdot 2^x$  arises. We have  $\lambda_1-\lambda_{19}$ : .0007; .0000; -.0040; -.0125; -.0196; -.0200; -.0290; -.0299; -.0317; -.0416; -.0416; -.0416; -.0458; -.0816; -.0874; -.0874; -.1181; -.1539; -.2222;  $\Lambda=-1.0672$ .

For  $k_8 = 0.075$  – regular attractor  $3 \cdot 2^0$  (see the window of periodicity in Fig. 2,b).  $\lambda_1 - \lambda_{19}$ : .0000; -.0004; -.0040; -.0117; -.0194; -.0211; -.0285; -.0285; -.0326; -.0406; -.0406; -.0406; -.0451; -.0819; -.0883; -.0883; -.1182; -.1563; -.2241;  $\Lambda = -1.0702$ .

For  $k_8=0.12$  – strange attractor  $2\cdot 2^x$ ;  $\lambda_1-\lambda_{19}$ : .0008; .0000; -.0040; -.0125; -.0192; -.0210; -.0287; -.0300; -.0324; -.0406; -.0406; -.0457; -.0822; -.0879; -.1172; -.1542; -.2212;  $\Lambda=-1.0653$ .

For  $k_8 = 0.27$  – strange attractor  $2 \cdot 2^x$ ;  $\lambda_1 - \lambda_{19}$ : .0002; .0000; -.0040; -.0125; -.0192; -.0219; -.0281; -.0310; -.0320; -.0387; -.0387; -.0387; -.0436; -.0842; -.0889; -.1185; -.1541; -.2209;  $\Lambda = -1.0638$ .

For  $k_8 = 0.275$  – strange attractor  $2 \cdot 2^x$ ;  $\lambda_1 - \lambda_{19} : .0001$ ; .0000; -.0040; -.0125; -.0192; -.0219; -.0280; -.0312; -.0323; -.0383; -.0383; -.0434; -.0842; -.0890; -.0890; -.1184; -.1540; -.2212;  $\Lambda = -1.0631$ .

For  $k_8 = 0.278$  – regular attractor  $2 \cdot 2^8$ ;  $\lambda_1 - \lambda_{19}$ : .0000; .0000; -.0041; -.0123; -.0193; -.0219; -.0283; -.0308; -.0320; -.0384; -.0384; -.0384; -.0435; -.0842; -.0890; -.1187; -.1538; -.2212;  $\Lambda = -1.0634$ .

For  $k_8=0.278$  – regular attractor  $2\cdot 2^8$ ;  $\lambda_1-\lambda_{19}$ : .0000; .0000; -.0041; -.0123; -.0193; -.0219; -.0283; -.0308; -.0320; -.0384; -.0384; -.0384; -.0435; -.0842; -.0890; -.1187; -.1538; -.2212;  $\Lambda=-1.0634$ .

For  $k_8 = 0.28$  – regular attractor  $2 \cdot 2^4$ ;  $\lambda_1 - \lambda_{19}$ : .0000; .0000; -.0041; -.0123; -.0193; -.0219; -.0282; -.0307; -.0322; -.0384; -.0384; -.0384; -.0436; -.0842; -.0890; -.1186; -.1540; -.2212;  $\Lambda = -1.0633$ .

The presented results of calculations indicate that the sum  $\Lambda$  of all indices, which determine the flow divergencies and, hence, the evolution of the phase volume along the trajectory, is maximal for the regular attractor  $3 \cdot 2^0$ . It arises in the window of periodicity for  $k_8 = 0.075$  ( $\Lambda = -1.0702$ ). For the strange attractors on the left and on the right (for  $k_8 = 0.01$  and  $k_8 = 0.12$ ), the divergencies are, respectively,  $\Lambda = -1.0672$  and  $\Lambda = -1.0653$ . This means that the phase volume element for the given attractor is contracted, on the whole, stronger along the trajectory. Here, we observe the self-organization of a stable cycle from chaotic modes. The Krebs cycle is adapted to the varying conditions.

By the given Lyapunov indices for strange attractors, we determine the KS-entropy (the Kolmogorov--Sinai entropy) [40]. By the Pesin theorem [41], the KS-entropy h corresponds the sum of all positive Lyapunov characteristic indices:

The KS-entropy allows us to judge about the rate, with which the information about the initial state of the system is lost. The positivity of the given entropy is a criterion of the chaos. This gives possibility to qualitatively estimate the properties of attractor's local stability.

We determine also the quantity inverse to the KS-entropy,  $t_{\min}$ . This is the time of a mixing in the system. It characterizes the rate, with which the initial conditions will be forgotten. For  $t << t_{\min}$ , the behavior of the system can be predicted with sufficient accuracy. For  $t > t_{\min}$ , only a probabilistic description is possible. The chaotic mode is not predictable due to the loss of the memory of initial conditions. The quantity  $t_{\min}$  is called the Lyapunov index and characterizes the "predictability horizon" of a strange attractor.

In order to classify the geometric structure of strange attractors, we calculated the dimension of their fractality. The strange attractors are fractal sets and have the fractional Hausdorff-Besicovitch dimension. But its direct calculation is a very labor-consuming task possessing no standard algorithm. Therefore, as a quantitative measure of the fractality, we calculated the Lyapunov dimension of attractors by the Kaplan--Yorke formula [42, 43]:

$$D_{F_r} = m + \frac{\sum_{i=1}^{m} \lambda_i}{\left|\lambda_{m+1}\right|},\tag{11}$$

where m is the number of the first Lyapunov indices ordered by their decreasing. Their sum  $\sum_{i=1}^m \lambda_i \geq 0$ , and m+1 is the number of the first Lyapunov index, whose value  $\lambda_{m+1} < 0$ .

For the above-considered strange attractors  $\,2^{\circ}$  , we obtained the following indices.

For  $k_8=0.01$ : h=0.0007,  $t_{\min}=1428.6$ ,  $D_{F_r}=2.175$ . For  $k_8=0.12$ : h=0.0008,  $t_{\min}=1250$ ,  $D_{F_r}=2.2$ . For  $k_8=0.27$ : h=0.0002,  $t_{\min}=5000$ ,  $D_{F_r}=2.05$ . For  $k_8=0.275$ : h=0.0001,  $t_{\min}=10000$ ,  $D_{F_r}=2.025$ .

By these indices, we can judge about the difference of the given strange attractors.

### **Conclusions**

With the help of the mathematical model of the Krebs cycle, we have studied the dependence of the cyclicity of the metabolic process on the amount of a final product of the oxidation, i.e., on the amount of the formed carbon dioxide. The multiplicity of the cycle is doubled by the Feigenbaum scenario, until the aperiodic modes of strange attractors arise. From them as a result of the self-organization, the stable periodic modes appear. This means that the system is adapted to the varying conditions. We have calculated the full spectra of Lyapunov indices and the divergencies for various modes. For the strange attractors, we have determined the KS-entropies, "predictability horizons," and Lyapunov dimensions of attractors. The results obtained allow us to study the structural-functional connections of the cycle of tricarboxylic acids, their influence on the cyclicity of metabolic oscillations in a cell, and the physical laws of self-organization in it.

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